



# Effect of mechanical stress on different core loss components along orthogonal directions in electrical steels

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## Abstract

The paper deals with the characterization and modelling of the mechanical stress dependency of magnetic losses along two orthogonal directions in non-oriented electrical steels. Significant anisotropy effects are highlighted. Using the three-term loss-separation approach, the different loss components are computed at each stress level for a wide range of frequency. Stress dependence of the core losses can be described in terms of the hysteresis and excess loss components, classical losses being assumed to be constant as a function of stress. Variations of the model coefficients with stress along the two principal directions are discussed. Such a model can be used for computing the losses in finite element analysis of rotating electrical machines or T-joint of transformers.

**Keywords** Core loss · Loss-separation approach · Mechanical stress · Electrical steels

## 1 Introduction

The design of electrical machines for more compact and high-speed systems demands reduction in size and weight. On the other hand, these systems are very sensitive to mechanical stresses which can be inherited from manufacturing stage (cutting, stacking, shrink fitting, etc.) or appear in operating conditions (magnetic forces, centrifugal forces, etc.) [1, 2]. The manufacturing processes are usually associated with residual stresses within the material which can significantly affect the material behaviour, even far from

plastically deformed areas. Mechanical stresses appearing in use usually remain within the elasticity range and also affect strongly the material behaviour [2, 3].

Numerous experimental studies on the influence of mechanical stress on the magnetic properties are reported in the literature [4–8]. Some of them show the influence of mechanical stress on magnetic properties using the direct observation of magnetic domains in grain-oriented materials [5, 6]. Effects of mechanical stress in non-oriented materials were reported recently [8] based on similar interpretation as for grain-oriented materials. It shows that the compressive stress monotonically deteriorates magnetic properties as compression stress disfavours the magnetic easy axis along the rolling direction. On the other hand, tensile stress first improves the magnetic properties, and then, after a certain value of applied stress, magnetic properties continuously deteriorate. This can be explained from a microstructural interpretation since the applied stress initially favours the magnetic easy axis, and at higher stress values, demagnetizing field may be dominant [6]. Other results on the effect of tensile stress and plastic deformation in electrical steels have been reported in [7].

These effects of stress are observed in particular on magnetic permeability and magnetic losses in core materials of electrical machines [9]. Indeed, ferromagnetic materials exhibit a strongly coupled magneto-mechanical behaviour [10, 11]. Therefore, the coupled behaviour needs to be

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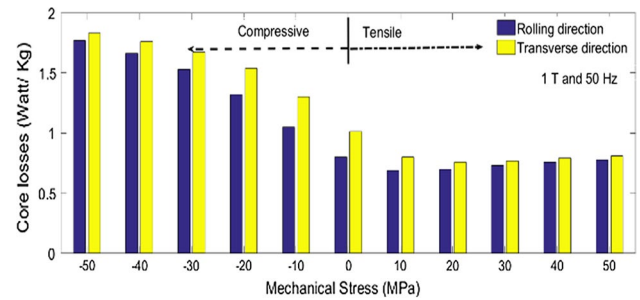
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taken into consideration in modern machine design tools. An accurate model for the description of magneto-mechanical effects is a prerequisite for such a tool. Numerous attempts have been reported in the literature which consider magneto-mechanical coupled behaviour into finite element (FEM) simulations [12–14]. However, the most common models are usually uncoupled and often consider uniform nonlinear anhysteretic magnetic behaviour in finite element analysis (FEA) simulations and iron losses to be computed at the post-processing stage [15–17]. This approach is also known as a posteriori approach. A significant number of papers on the magneto-elastic constitutive law for ferromagnetic materials can be found in the literature [18–22]. However, the numerical implementation of these models in FEM (also known as *intrinsic* approaches) is not a straightforward task as it takes huge computation time and poses convergence issues as well. Hence, the former is generally a preferred choice if its accuracy is within an acceptable range. For such an analysis, anhysteretic magnetic and iron loss behaviour should consider the influence of mechanical stresses.

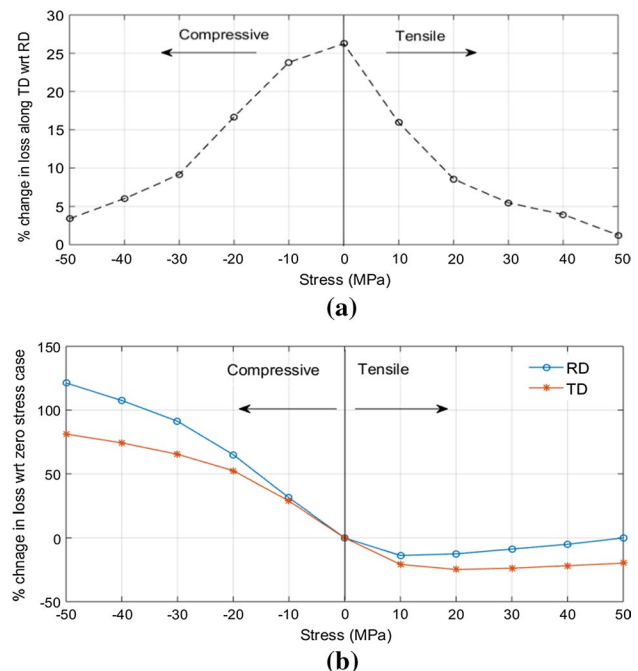
According to statistical loss theory, the core losses can be classified into three components, namely hysteresis, classical eddy current and excess losses [23]. A number of loss formula based on the loss-separation approach have been reported in the literature [24–27], which can be applied to compute core losses at post-processing stage in FEM analyses of electrical machines.

A significant number of studies are reported on the description of stress effects on the core losses using the loss-separation approach [28–32]. The behaviour of loss components under applied mechanical stress has been studied, and the loss coefficients modelled using the exponential functions [28] or second-order polynomial [30]. Although the effect of compressive and tensile stress on the iron losses has been studied extensively in the literature [9, 11, 28–32], most of the existing attempts of loss modelling are limited to the rolling direction (RD) only [28, 30, 32]. The effect of stress on magnetic losses in other than rolling direction is rarely addressed in the literature [3]. However, even if the anisotropy of non-oriented iron–silicon steels is known to be moderate, significant differences can be observed in the effect of stress on core losses along the rolling and transverse (TD) directions. This anisotropy is illustrated in Fig. 1 for a non-oriented material (M235-35A) at 1 T and 50 Hz subjected to uniaxial tensile and compressive stress.

The percentage difference in core losses along RD and TD at different stress levels (compression and tension) is depicted in Fig. 2a. The maximum difference between the losses in the two orthogonal directions occurs at zero stress condition and decreases as mechanical stress increases [3]. The application of stress then tends to uniformize the behaviour and hence reduce the anisotropy. This can be connected to the modification in the domain configuration which



**Fig. 1** Core losses along two principal directions (rolling—RD, and transverse—TD) under applied uniaxial mechanical stress. The stress is applied along the magnetic field direction



**Fig. 2** Effect of mechanical stress on the core losses along two principal directions at 1 T, 50 Hz **a** % change in loss along TD with respect to RD at different stress levels, **b** % change in loss with respect to unstressed condition along RD and TD

becomes more and more dominated by stress effects as the level of stress increases [8]. The difference in the core loss with respect to zero stress condition is shown in Fig. 2b. A significant change can be observed in both principal directions as the applied stress increases. It is evident that compression significantly increases the magnetic losses, while tension has a more moderate, and initially beneficial, effect.

Loss properties being one of the key-points for design and development of high power density and efficient machines, there is a need for the development of loss models considering the effects of stress. Moreover, the magnetic flux does not remain along a single direction, especially, in joint and teeth regions of electrical machines. Therefore,

the anisotropic feature of the core losses in electrical steels should be considered for an accurate loss computation in electrical machines [33]. This is why it is useful to describe loss behaviour in electrical steels along different directions. This can help designers and researchers in the development of efficient loss models. Furthermore, studies on magneto-mechanical effects are vital for optimization of electrical machines, determination of lifetime performance and material choice.

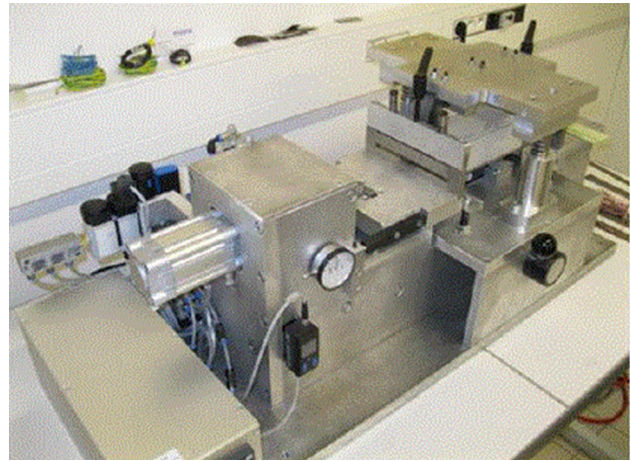
This paper investigates the effect of stress on the core losses in two orthogonal directions. The paper also presents an approach to model these effects in both principal directions. The suggested approach is based on the loss separation and investigates the effects of mechanical stress using a frequency-domain loss model. As evident from Fig. 1, compressive stresses affect the core loss more significantly than the tensile stress. Therefore, this paper focuses on the modelling of compressive stress effects on the core losses along the two principal directions. The paper also discusses variations of loss components with stress. The loss-separation approach is discussed in Sect. 2, and mechanical stress dependency of loss model parameters is given in Sect. 3 for both rolling (RD) and transverse (TD) directions.

## 2 Measurements and modelling of core losses

### 2.1 Experimental set-up

The measurements of hysteresis loops and corresponding losses are taken on samples (non-oriented electrical steel; grade-M235-35A) which are cut at two angles of  $0^\circ$  and  $90^\circ$  with respect to RD. The measurements are taken on a single sheet tester (SST), with tension and compression mechanical unit as shown in Fig. 3. The magnetic measurement system (SST) is a BROCKHAUS-MPG 200 D device. The measurements are made under standard pure sinusoidal waveforms of induction levels (0.2–1.5 T) for a frequency range of 10 Hz–2 kHz. Each measurement is taken twice, and the results are averaged. Stress is applied through a pneumatic tension and compression unit. Given the dimensions of the sample (standard stripe, 30 mm width and 0.35 thickness), the pneumatic unit is capable of applying uniaxial stresses up to 70 MPa. The resolution on stress is approximately 3.5 MPa. During compression, a wedge is placed on each side of the sample in order to avoid buckling. This allows loading the sample up to 50 MPa in compression.

From the measured B–H curves, total core losses at each frequency and stress level are calculated by computing the area of hysteresis loops. The separation of total core loss into



**Fig. 3** Single sheet tester with the mechanical unit for stresses (Source: RENAULT-SAS, Guyancourt, France)

three components, viz. hysteretic, classic and excess terms, is a common practice for loss description in thin ferromagnetic laminations [34].

### 2.2 Loss-separation model

The loss mechanisms associated with the magnetization process are a complex phenomenon that is still far from being understood [35]. A standard physics-based loss computational model is yet to achieve. Empirical approaches can be used to describe iron losses in an approximate manner [36]. The core losses, from an engineering viewpoint, can be separated in three components, namely the hysteresis, classical eddy current and excess or anomalous losses [9, 34]. In recent years, a number of advances have been proposed in the loss-separation models [25] in order to consider harmonics or minor loops [15, 16, 27].

For electrical steels, the decomposition of the total loss into three components is a common practice. The total loss can be rewritten as [34]:

$$P_{\text{tot}} = P_{\text{h}} + P_{\text{ec}} + P_{\text{ex}} \quad (1)$$

where  $P_{\text{tot}}$  is the total core loss (W/kg),  $P_{\text{h}}$  is the hysteresis loss,  $P_{\text{ec}}$  is the eddy current loss, and  $P_{\text{ex}}$  is the excess or anomalous losses.

The static hysteresis losses can be explained through the irreversible motion of domain walls [37]. The losses at very low frequency can be referred to as the quasi-static or DC hysteresis loss which is computed by the area of DC hysteresis loops. This component is strongly related to the domain micro-structure of materials [35]. It can be represented as:

$$P_{\text{h}} = k_{\text{h}} \hat{B}^{\alpha} f \quad (2)$$

where  $k_h$  and  $\alpha$  are the material constants which depend on the domain structure and crystalline orientations of materials.

The classical eddy current loss depends on the supply frequency, sample thickness and conductivity of laminations under the assumption of homogeneity of the material [34]. Conductivity of the sample is usually assumed to be isotropic and independent of stress. The eddy current losses can be calculated using the following equation:

$$P_{ec} = k_{ec} \hat{B}^2 f^2 \quad (3)$$

where  $k_{ec}$  is a coefficient which depends on the geometrical dimensions and the chemical composition of the material.  $\hat{B}$  and  $f$  are the peak induction value and frequency, respectively.

An excess power loss that depends on domain configuration and microscopic eddy currents is modelled as [35]:

$$P_{exc} = k_{ex} \hat{B}^{1.5} f^{1.5} \quad (4)$$

where  $k_{ex}$  is a material coefficient and it depends on the domain structure in the material.

According to conventional models, the eddy current loss, also referred to as the classical loss, can be calculated directly as [24]:

$$k_{ec} = \frac{(\pi d)^2}{6\rho D} \quad (5)$$

where  $d$ ,  $\rho$  and  $D$  are the sheet thickness ( $m$ ), resistivity ( $\Omega m$ ) and mass density ( $kg/m^3$ ) of the core material, respectively. Due to inherent limitations of conventional models (with four coefficients to identify), the accuracy of loss computation may not be within acceptable limits for a wide range of frequencies and induction levels. On contrary, an effective loss model is essential for accurate design of electrical machines, where varying magnetic loadings and flux weakening at high speeds in variable speed machines are often encountered [25]. Therefore, numerous empirical models have been proposed for loss evaluation over wide ranges of frequencies and induction levels [25, 36].

A three-term loss model using a modified Steinmetz formula has been given as [36]:

$$P_{tot} = k_h \hat{B}^\alpha f + k_{ec} \hat{B}^2 f^2 + k_{ex} \hat{B}^{1.5} f^{1.5} \quad (6)$$

In this approach, the measured loss data of a particular material over a range of frequency and induction levels are fitted to a loss prediction formula and the unknown parameters of the model are determined based on a data fitting algorithm as discussed in [24]. Moreover, the loss model (6) has been

further modified in order to consider dynamic hysteresis and skin effects at higher frequencies as [38]:

$$P_{tot} = k_h \hat{B}^{(\alpha_0 + \alpha_1 B + \alpha_2 B^2)} f + k_{ec} \hat{B}^2 f^2 + k_{ex} \hat{B}^{1.5} f^{1.5} \quad (7)$$

Here, the parameters ( $k_{ec}$ ,  $k_{ex}$ ) vary with induction levels and can be represented as a function of  $B$  [36]. However, hysteresis loss coefficients ( $k_h$ ,  $\alpha$ ) are frequency-dependent. Moreover, the hysteresis exponent coefficient ( $\alpha$ ) is represented as a second-order polynomial of  $B$  which considers the skin effect and dynamic hysteresis [38]. The algorithm, however, is simple and valid for a frequency range from static up to 2 kHz [36].

### 2.3 Loss coefficient identification

The loss coefficients ( $k_{ec}$ ,  $k_{ex}$ ,  $k_h$  and  $\alpha$ ) can be obtained using the algorithms presented in [24, 36]. By dividing the total losses by the frequency, Eq. (7) can be rewritten as:

$$P_{tot}/f = a + bf + c\sqrt{f} \quad (8)$$

Here, the coefficients  $a$ ,  $b$  and  $c$  are represented as:

$$\begin{aligned} a &= k_h \hat{B}^\alpha \\ b &= k_{ec} \hat{B}^2 \\ c &= k_{ex} \hat{B}^{1.5} \end{aligned} \quad (9)$$

The measured loss values at least for three frequencies are required for determining three coefficients  $a$ ,  $b$  and  $c$  in (8) at any induction level. From (8) and (9), the eddy current ( $k_{ec}$ ) and excess loss coefficients ( $k_{ex}$ ) can easily be computed. These coefficients show a significant variation with induction levels but no dependence on frequency [24]. A second-order polynomial is used to fit these variations as:

$$\begin{aligned} k_{ec} &= k_{ec0} + k_{ec1} B + k_{ec2} B^2 \\ k_{ex} &= k_{ex0} + k_{ex1} B + k_{ex2} B^2 \end{aligned} \quad (10)$$

In order to identify the hysteresis loss coefficients ( $k_h$ ,  $\alpha$ ), Eq. (8) can be rewritten using the logarithmic operator as:

$$\log a = \log k_h + (\alpha_0 + \alpha_1 \cdot B) \cdot \log B \quad (11)$$

The hysteresis exponent coefficient ( $\alpha$ ) is taken in this work as a linear function of  $B$  in order to reduce the unknown parameters and required measured data for parameter identification. For three unknowns, Eq. (11) needs the loss data at least for three induction levels. The coefficient  $a$  represents the ratio of hysteresis loss and frequency, and it can be computed using (8) by substituting the values of  $b$  and  $c$  which are obtained from (9).



**Table 1** Determined loss coefficients

Parameters	Identified values (unit less)
$k_h$	0.0121
$\alpha (\alpha_0, \alpha_1)$	$\alpha_0 = 1.6829; \alpha_1 = 0.4050$
$k_{ec} (k_{ec0}, k_{ec1}, k_{ec2})$	$k_{ec0} = 3.9305 \times 10^{-5};$ $k_{ec1} = -4.4837 \times 10^{-4};$ $k_{ec2} = 3.9840 \times 10^{-5}$
$k_{ex} (k_{ex0}, k_{ex1}, k_{ex2})$	$k_{ex0} = 6.2871 \times 10^{-5};$ $k_{ex1} = 0.0013;$ $k_{ex2} = -8.4956 \times 10^{-4}$

### 2.4 Modelling of core losses using the loss-separation theory

The loss model coefficients under no applied stress are identified using the measured loss data at three induction levels (here 0.2, 0.7 and 1.5 T) and three frequencies (here 50, 200 and 800 Hz) for the frequency range 10 Hz–1 kHz. The obtained loss coefficients are given in Table 1.

The computed and measured losses are shown to be in close agreement at various induction levels and for a wide frequency range as depicted in Fig. 4a–c. The magnitude of maximum error for a frequency has not crossed beyond 12% in this analysis as shown in Fig. 4b.

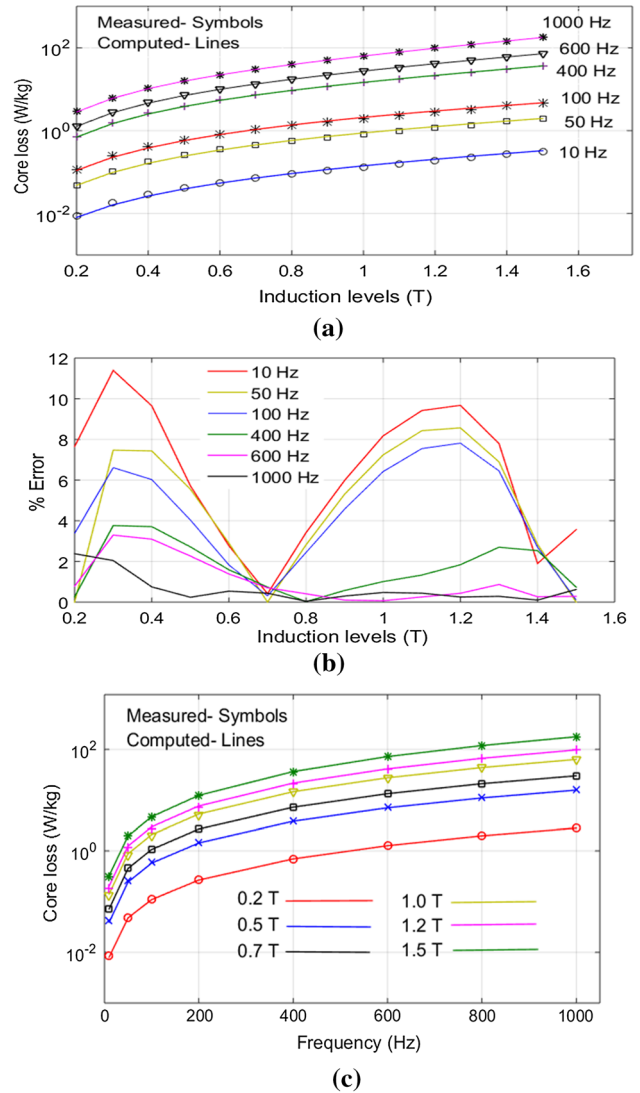
The eddy current and excess loss parameters can then be represented as a second-degree polynomial function of induction ( $B$ ; as given in Eq. 10) to get better accuracy over a wide range of induction values. Variations of these parameters ( $k_{ec}$ ,  $k_{ex}$ ) and different loss components with induction levels are shown in Fig. 5a, b. Variations of the two loss coefficients ( $k_{ec}$ ,  $k_{ex}$ ) with induction levels are complementary-like to each other. Moreover, the excess loss is the maximum between 0.6 and 0.8 T and then decreases at higher induction levels. Eddy current losses are minimum between 0.6 and 0.8 T and then start to increase at higher induction values.

It is shown that the suggested model (7) can satisfactorily describe the loss behaviour of this material. It is used in the next section to study the effect of mechanical stresses on the core losses in the same material.

## 3 Investigation of loss behaviour along the rolling and transverse directions

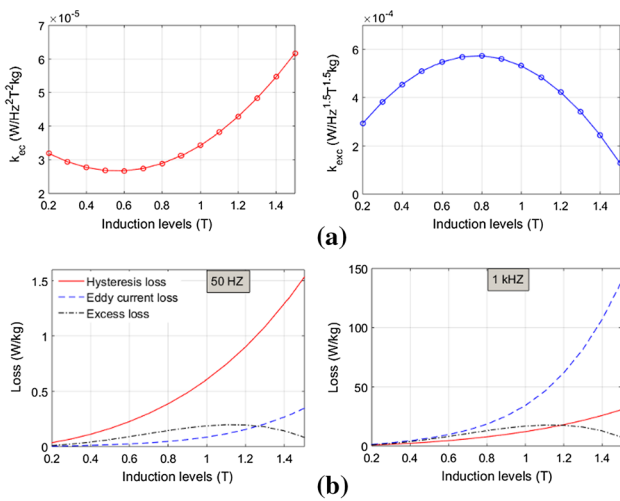
### 3.1 Loss–stress behaviour along the rolling and transverse directions

In order to investigate the stress effects on the losses along the two orthogonal directions, the loss-separation approach is used to compute each component along these directions

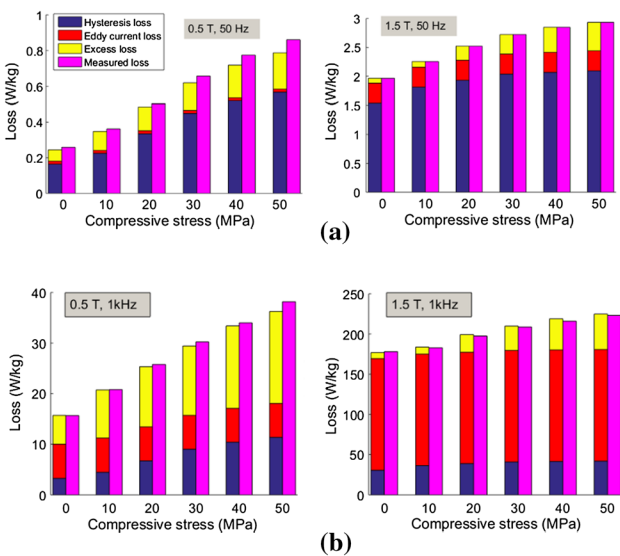


**Fig. 4** Computed and measured core loss: **a** losses versus induction, **b** percentage error versus induction levels at different frequencies and **c** losses versus frequency (the loss values at 0.2, 0.7 and 1.5 T and at 50 Hz, 200 Hz and 800 Hz are used in the identification algorithm)

at different stress levels. The loss separation is performed along RD and TD to investigate the behaviour of three loss components at different stress levels. The classical loss component is assumed to be independent of stress and orientation. Hence, this component is assumed to remain the same for all directions and stress levels and it can be computed from RD measurement at zero stress level. The hysteresis and excess loss components vary with the mechanical stress as they depend on the domain configuration of the material. Therefore, two loss coefficients need to be identified (hysteresis and excess loss coefficient). They can be determined using the measured loss data [at three induction levels (here 0.2 T, 0.7 T and 1.5 T) but only for two frequencies (here 50 and 800 Hz)] for each stress.



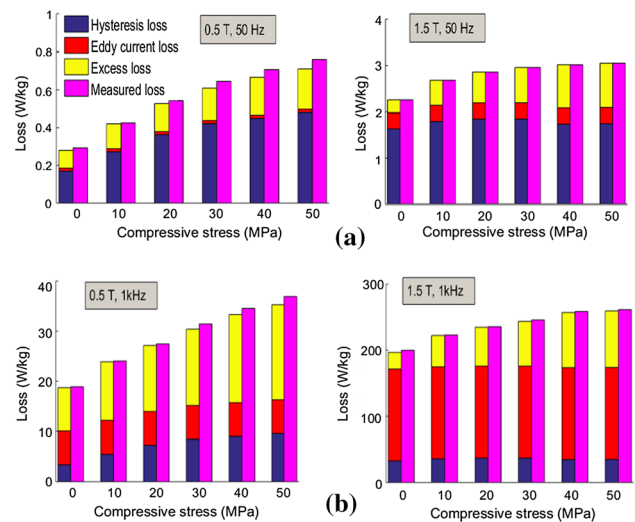
**Fig. 5** Variation of **a** loss coefficients with induction levels and **b** loss components with induction levels at different frequencies



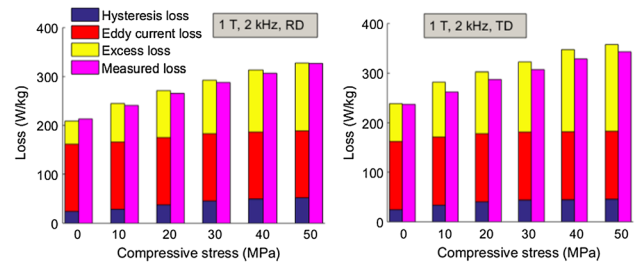
**Fig. 6** Computed loss components and measured core losses with compressive stress along RD

The computed loss components and the measured core loss are in reasonably close agreement at different induction levels and frequencies, as shown in Fig. 6a, b.

The same approach is carried out on the core loss data along TD. It is reminded here that the eddy current loss will remain the same at each stress level and whatever the direction of the applied magnetic field. Therefore, parameter identification is performed at the same induction levels but only for two frequencies (here 50 Hz and 800 Hz) as used for RD at nonzero stress levels. Computed and measured losses along TD at different induction levels and frequencies are shown in Fig. 7a, b.



**Fig. 7** Computed loss components and measured core losses with compressive stress along TD



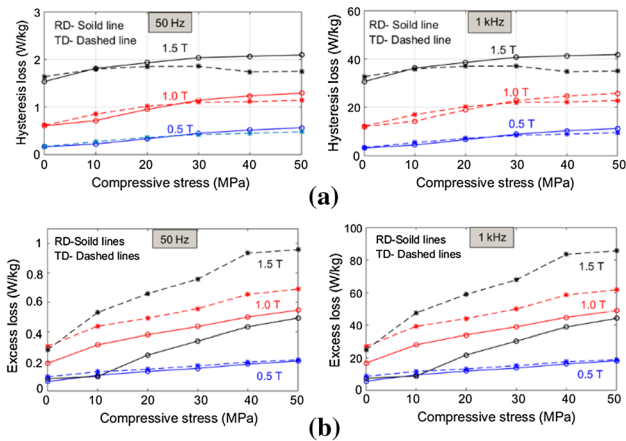
**Fig. 8** Predicted and measured losses along RD and TD

Using the identified parameters, the model can be used to predict core losses at other induction levels and frequency (e.g. 1.0 T, 2 kHz) with reasonable accuracy as shown in Fig. 8. The magnitude of maximum error for a frequency has not crossed beyond 15% in this analysis. The accuracy of the analysis can be further improved by using higher-order polynomial for  $k_{ec}(B)$ ,  $k_{ex}(B)$  and  $\alpha(B)$  as discussed in [25].

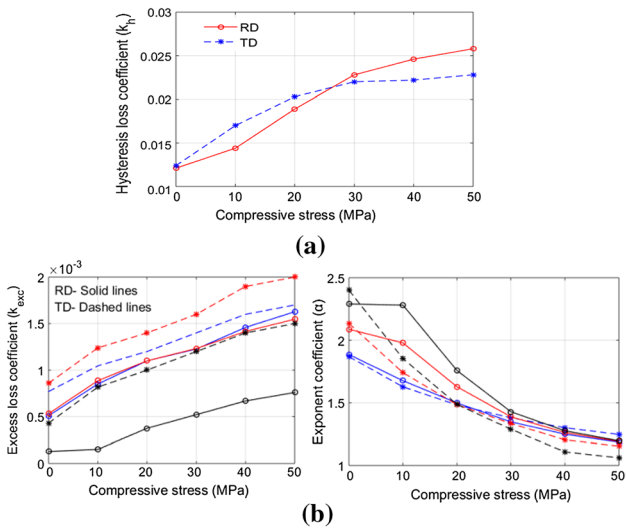
### 3.2 Results and discussion

The effects of microstructural features (crystallographic texture, grain size, dislocation densities, residual stresses, etc.) on the core losses can be described in terms of  $k_{ex}$  and  $k_h$  [35]. These parameters are a signature of the effect of mechanical stress on the microstructure and magnetic properties of electrical steels [29]. Hysteresis and excess loss variations with compressive stress show similar trends as shown in Fig. 9a, b.

It can be observed from Fig. 9 that both hysteresis and excess components show similar trends (monotonically increasing with stress) at different frequency and induction



**Fig. 9** Variation of loss components with compressive stress at different frequencies and induction levels (blue)=0.5 T; red=1.0 T; black=1.5 T; circle (solid lines)=RD; and cross (dashed lines)=TD **a** hysteresis component ( $P_{hyst}$ ) and **b** excess loss component ( $P_{exc}$ ) (colour figure online)



**Fig. 10** Variation in loss coefficient with compressive stress along RD and TD **a**  $k_h$  and **b**  $k_{ex}$  and  $\alpha$  (blue)=0.5 T; red=1.0 T; black=1.5 T; circle (solid lines)=RD; and cross (dashed lines)=TD (colour figure online)

levels. The excess loss component for TD, as shown in Fig. 9b, is higher than that of RD up to 50 MPa. On the other hand, the hysteresis loss component along TD is higher than that of RD up to a certain stress level and beyond this stress level, it becomes lower than RD. The obtained results are in line with those reported in [3] which uses although a two-term loss formula. In order to investigate the stress dependency of loss components, variations of loss coefficients with stress are shown in Fig. 10a, b. Parameters for hysteresis ( $k_h$ ) and excess ( $k_{ex}$ ) losses increase monotonically with the compression stress level. On the other hand, the hysteresis exponent coefficient ( $\alpha$ ) decreases as the applied stress increases.

These parameter variations with stress have been modelled using exponential function [28] and second-order polynomial [30] in the literature. However, these works are restricted to RD. The present work provides an approach to model stress dependencies on loss parameters along both principal directions. The difference in the loss components for the two orthogonal directions might be attributed to the anisotropy due to the crystallographic texture of the material. That is the reason why only two components (hysteresis and excess loss which depend on domain configurations) are shown to vary with respect to stress and direction as well. On the other hand, the classical eddy current loss is assumed to be constant with respect to applied stress and orientation of magnetic field. This is consistent with previous works [28, 30].

Variations of model parameters with stress show almost similar trends along the two directions. Therefore, it can be concluded that the loss–stress dependency shows similar trends along both the principal directions. In other words, the stress can be included in the suggested loss model representing its coefficients as a similar function of stress along both directions. The accuracy of the loss model when stress effects are included can be improved by choosing appropriate stress-dependent functions for the loss coefficients [28, 30]. Of course, such a single-dimensional approach has to be generalized to multiaxial configurations, either by using multiaxial measurement apparatus [39–41] or based on 3D modelling approaches [42, 43].

### 4 Conclusions

The effect of uniaxial mechanical stress, particularly compressive stress, on iron losses is noticeable in both principal directions. The losses under 50 MPa compression stress can be 120% higher than that under unstressed condition along RD and 85% higher along TD. Therefore, these effects should be considered in loss computations for electrical machines. It can also be observed that at higher stress levels, the effect of stress is highly detrimental to the loss performance and the difference between RD and TD vanishes. In this paper, using a modified empirical loss model, stress dependencies of the hysteresis and excess loss components along RD and TD are studied for a wide range of frequencies (from static up to 2 kHz) and induction levels (0.1–1.5 T). The eddy current loss component is assumed to be independent of mechanical stress. The variations of loss coefficients with compressive stress are investigated in the two principal directions. These variations in hysteresis and excess loss parameters against the stress level are almost similar along the two directions. It should be noted that the proposed analysis is valid only for the case of elastic mechanical loadings. The stress–loss behaviour does not show strong orientation dependencies as the stress dependencies of different loss

components follow the same trends along both direction. Higher losses are nevertheless observed along TD compared to RD which is attributed to the anisotropy of the material resulting from the crystallographic texture.

The computed core losses using the suggested loss model match with the measured total core losses along RD and TD. Such a study can be useful for a better prediction of the iron losses in rotating electrical machines [33]. Implementation of the model in FEM analysis of electrical machines is identified as a future work.

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